RESEARCH

In-flight Scalar Calibration and Characterisation of the Swarm Magnetometry Package

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Abstract

We present the in-flight scalar calibration and characterisation of the *Swarm* magnetometry package consisting of the absolute scalar magnetometer (ASM), the vector magnetometer (VFM), and the spacecraft structure supporting the instruments. A significant improvement in the scalar residuals between the pairs of magnetometers is demonstrated, confirming the high performance of these instruments. The results presented here, including the characterization of a Sun-driven disturbance field, form the basis of the correction of the magnetic vector measurements from *Swarm* which is applied to the *Swarm* Level 1b magnetic data.

Keywords: Geomagnetism; Magnetometer; Instrument Calibration; Satellite; *Swarm*

1 Introduction

In November 2013 the European Space Agency (ESA) launched the three Swarm satellites named Alpha, Bravo, and Charlie with the objective to provide the best ever survey of the geomagnetic field and its temporal evolution (Friis-Christensen et al., 2006). Each spacecraft carries an Absolute Scalar Magnetometer (ASM) for measuring Earth's magnetic field intensity, a Vector Fluxgate Magnetometer 8 (VFM) measuring the direction and strength of the magnetic field, and a three-head 9 Star TRacker (STR) mounted close to the VFM to obtain the attitude needed to 10 transform the vector readings to an Earth fixed coordinate frame. Time and position 11 are provided by an on-board GPS receiver. The payload also includes instruments 12 to measure plasma and electric field parameters as well as non-gravitational ac-13 celeration. More information on the mission status after two years in orbit can be 14 found in *Floberghagen et al.* (2016). 15

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One of the purposes of the scalar magnetometer (ASM) is to provide the nec-17 essary absolute magnetic data to calibrate the vector magnetometer (VFM). For 18 this an approach similar to that adopted for the previous satellite missions Ørsted 19 and CHAMP was foreseen (c.f. Olsen et al., 2003; Yin and Lühr, 2011) since those 20 missions carried equivalent instrumentation. However, soon after launch of Swarm 21 it became clear that the magnetic field vector measurements on all three space-22 craft were contaminated by unforeseen disturbances which could not be captured 23 by the traditional in-flight calibration methods referred to above. Furthermore, the 24 disturbances show systematic variation which could impact or map into scientific 25 investigations based on *Swarm* magnetic data. The light blue symbols in Fig. 1 show

time series of the *scalar residuals*, which are the difference, $\Delta F = |\vec{B}_{\rm VFM}| - F_{\rm ASM}$, 27 between the modulus of the VFM data, $|\vec{B}_{\rm VFM}|$, and the magnetic intensity mea-28 surements, F_{ASM} , taken by the ASM instrument. Based on experience with Ørsted 29 and CHAMP scalar residuals with sub-nanotesla level were expected (rms value 30 well below 0.5 nT), while for *Swarm* the scatter of the residuals was observed to 31 reach several nT, resulting in an rms value approaching 1 nT, but crucially show-32 ing a very clear Local Time dependence. A task force was therefore established to 33 investigate and mitigate the effect. 34

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Detailed investigations of the scalar residuals ΔF and of the ASM and VFM measurements separately indicated that:

• the vector readings of the VFM are affected by a disturbance vector field;

• the scalar readings of the ASM are much less, if at all, affected.

Consequently the Task Force concluded to pursue models which assume the mag-40 netic disturbance to be affecting the VFM measurements only. Plotting ΔF as a 41 function of the Sun incidence angles with respect to the spacecraft, reveals system-42 atic features of the disturbance, as shown in Fig. 3. At the start of section 2 we 43 provide detailed definitions of the two Sun incidence angles α and β . This supports 44 the hypothesis that a magnetic source in the vicinity of the VFM magnetometer, 45 with strength and direction depending on the direction to the Sun (as seen from the 46 spacecraft), is responsible. We refer to such a disturbance field vector, that depends 47 on the direction to the Sun, as $\delta \vec{B}_{Sun}$. 48

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The purpose of this article is to document the details of in-flight calibration of the *Swarm* magnetometer package, including an empirical determination and removal of the Sun driven vector disturbance field $\delta \vec{B}_{Sun}$, based on a mitigation approach proposed by Vincent Lesur (*Lesur et al.*, 2015).

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Section 2 describes the parameterisation of the model of the Sun-driven distur-55 bance – in following referred to as the *characterisation* of the disturbance field 56 and of the *calibration* of the VFM instrument, by which means determination 57 of its intrinsic scale factors and their dependence on time and temperature, and 58 determination of the sensor-axis non-orthogonalities. We document the adopted 59 Iteratively Reweighted Least Squared (IRLS) estimation approach, that includes 60 a truncated singular value decomposition (SVD) approach to solving the inverse 61 problem. The results obtained for Swarm Alpha, based on data covering the pe-62 riod from launch (22. November 2013) until end of June 2015 (i.e. 19 months), are 63 presented in Section 3. Application of the scheme to data from the satellites Bravo 64 and *Charlie* resulted in similar levels of residual improvement and statistics, and 65 the estimates of the Sun driven disturbance $\delta \vec{B}_{Sun}$ show generally similar behaviour 66 and structural features as found for Swarm Alpha, although there are also some 67 differences. Finally, Section 4 summarizes the findings and provides perspectives 68 regarding further improvements of the method. 69

71 2 Characterisation and Calibration with Scalar Residuals

The Sun incidence angles α and β are crucial in our approach to *characterise* the 72 scalar residual. To clarify, in Fig. 2 we illustrate the definition of these angles with 73 respect to the spacecraft and the Sun position. α is the azimuth in the spacecraft 74 x-z plane (nominally the orbit plane) and β is the "elevation" out of the x-z plane 75 positive towards *left* (looking in the nominal flight direction; i.e. positive *opposite* 76 the spacecraft y axis). Examples of values for α and β for particular Sun positions 77 are: Considering how these angles vary over orbits of the Swarm spacecraft, we find 78 the angle α varies rapidly: from 0° to 360° within one orbit (i.e. within \approx 90 minutes) 79 while the other angle, β , varies more slowly (by $\approx 1.5^{\circ}$ in one day). 80

- $\beta = +90^{\circ}$: Sun directly from -y (i.e. from the left during nominal flight)
- $\beta = -90^{\circ}$: Sun directly from +y (i.e. from the right)
- 82 83

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• $\beta = 0^{\circ}, \alpha = +90^{\circ}$: Sun directly from -z (above)

• $\beta = 0^{\circ}, \alpha = 0^{\circ}$: Sun directly from +x (i.e. from the front)

84 85 86

• $\beta = 0^{\circ}, \alpha = +180^{\circ}$: Sun directly from -x (i.e. from the back – slightly above the boom)

⁸⁷ Considering how these angles vary over orbits of the *Swarm* spacecraft during nom-⁸⁸ inal flight, we find that α varies rapidly: from 360° down to 0° within one orbit (i.e. ⁸⁹ within \approx 90 minutes) while β , varies slowly up and down typically by $\approx 1.25^{\circ}$ in ⁹⁰ one day (for *Alpha* and *Charlie*, 1.20° for *Bravo*).

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Although the observed scalar residuals clearly vary with the Sun incidence angles 92 α and β (see Fig. 3) there is no direct mapping of ΔF in terms of these parameters. 93 This is a consequence of the scalar residuals $\Delta F \approx \delta \vec{B}_{Sun} \cdot \vec{b}_0$ being the projection 94 of the magnetic disturbance vector $\delta \vec{B}_{Sun}$, onto the unit vector \vec{b}_0 of the ambient 95 magnetic field direction (Earth's main field). The former is oriented relative to 96 the spacecraft while the latter is oriented relative to Earth, which results in the 97 variations with the spacecraft local time (captured by β) as seen in Fig. 3. The 98 spacecraft local time changes by 12 hours (corresponding to a change in β by 180°) 99 within approximately $4\frac{1}{2}$ months. 100

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To account for the projection on to the ambient field, we consider a *vector* magnetic disturbance $\delta \vec{B}_{Sun}(\alpha, \beta)$, with each component depending individually on the Sun incidence angles. Mathematically, we describe each component of the disturbance field vector by a spherical harmonic expansion in α and β i.e. we consider three independent spherical harmonic expansions in all.

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This model *characterizing* the Sun-driven disturbance is co-estimated together 108 with a model of the temporal evolution of the VFM sensitivity and an adjustment 109 of the pre-flight estimated non-orthogonality angles of the VFM sensor. For this 110 we perform a *scalar calibration* via a least squares fit, minimizing the discrepancy 111 (ΔF) between the fully calibrated and corrected measurements from the ASM and 112 the modulus of the vector measurements from the VFM after our model has been 113 applied. Huber weights are used iteratively to eliminate the effect of anomalous 114 measurements ("outliers") on the estimated models. 115

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117 2.1 Model Parameterisation

As outlined above, our model characterizing the Sun-driven disturbance vector $\delta \vec{B}_{Sun}$ consists of three spherical harmonic expansions up to degree and order 25, one for each of the magnetic field components in the VFM magnetometer frame, with the position of the Sun with respect to the spacecraft parameterised by the Sun incidence angles α and β . It takes the form

$$\delta \vec{B}_{\rm Sun} = \sum_{n=0}^{25} \sum_{m=0}^{n} \left(\vec{u}_n^m \cos m\alpha + \vec{v}_n^m \sin m\alpha \right) P_n^m(\sin \beta)$$

where \vec{u}_n^m and \vec{v}_n^m are the spherical harmonic expansion coefficients, with one component for each component of the disturbance field, and P_n^m are the Schmidt seminormalized Legendre functions. Note that $\delta \vec{B}_{Sun}$ includes static terms (n = m = 0), that describe a static (i.e. independent of the Sun position) disturbance vector. The disturbance field vector $\delta \vec{B}_{Sun}$ is thus described by $3 \times 26^2 = 2,028$ model coefficients.

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The model for re-scaling the vector measurements and taking into account any small adjustment of the non-orthogonality of the VFM sensors, which is required in order to obtain the fully calibrated and corrected vector field measurements $\vec{B}_{\rm VFM}$, now takes the form

$$\vec{B}_{\rm VFM} = \underline{P}^{-1} \underline{S}^{-1} \vec{B}_{\rm pre-flight} - \delta \vec{B}_{\rm Sur}$$

where $\vec{B}_{\rm pre-flight}$ are the VFM measurements calibrated using the pre-flight parameters and corrected for the pre-flight determined stray fields as described in $T \emptyset ffner-$ *Clausen* (2015). <u>S</u> is a 3 × 3 diagonal scaling matrix with elements

$$s_j = s^{\mathrm{B-spline}}(t) + s_{j,\mathrm{Tsensor}}T_{\mathrm{sensor}} + s_{j,\beta}\beta$$

where $s^{B-spline}(t)$ is a quadratic B-spline in time with 3-month knot separation 137 (common for all three components of the magnetic field), and $s_{j,\text{Tsensor}}, j = 1 - 3$ 138 is an adjustment of the pre-flight estimated dependency of the VFM sensitivity on 139 its sensor temperature, T_{sensor} , for each sensor axis j. $s_{j,\beta}$ is an empirical scaling 140 parameter and β the Sun incidence angle, as defined above. The choice of quadratic 141 B-splines with 3-month knot separation is made to allow sufficient flexibility of the 142 model; the exact choice of B-spline knot times is not crucial as very similar results 143 are obtained with other, similar parameterisations. The estimated B-splines exhibit 144 very moderate accelerations (in the case of the full model, see Fig. 6) and it may be 145 possible to simplify the parameterisation of the time-dependence in future models, 146 e.g. to an exponential saturation in time as this is the expected behaviour of the 147 VFM instrument sensitivity, however an exponential model is ill-conditioned on the 148 timespan of data used here. 149

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 $\underline{\underline{P}}$ is the non-orthogonality matrix that makes small adjustments to the pre-flight estimated non-orthogonalities of the VFM sensor (cf. *Olsen et al.*, 2003)

$$\underline{\underline{P}} = \begin{pmatrix} 1 & 0 & 0 \\ -\sin u_1 & \cos u_1 & 0 \\ \sin u_2 & \sin u_3 & \sqrt{1 - \sin^2 u_2 - \sin^2 u_3} \end{pmatrix}$$

Our in-flight calibration model comprises 18 parameters in all; together with the 2,028 parameters describing $\delta \vec{B}_{Sun}$ this results in 2,046 model parameters to be estimated, as listed in Table 1.

156 2.2 Estimation of Model Parameters: Inversion and Regularisation

In order to estimate the 2,046 model parameters from the scalar residuals we need
to solve a nonlinear inverse problem. The nonlinearity arises from the the treatment
of non-orthogonalities (*Olsen et al.*, 2003).

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The forward relationship between the vector of the scalar residuals, \mathbf{d} , $(d_i = \Delta F_i)$, the scalar residual of the *i*th data point) and the model parameter vector \mathbf{m} , may therefore be written in the form

$$\mathbf{d} = \mathbf{g}(\mathbf{m}) + \mathbf{e}$$

where $\mathbf{g}(\mathbf{m})$ is a nonlinear function of the models parameters and \mathbf{e} is a small remainder, that cannot be explained by the model, which we seek to minimise.

Linearisation of this problem is straightforward. A *regularized*, *iterativelyreweighted*, *least squares solution* to the inverse problem, is then obtained using the algorithm

$$\mathbf{m}_{k+1} = \mathbf{m}_k + (\underline{\mathbf{G}}_k^T \underline{\mathbf{W}}_k \underline{\mathbf{G}}_k + \lambda \underline{\mathbf{R}})^{-1} \left(\underline{\mathbf{G}}_k^T \underline{\mathbf{W}}_k \left[\mathbf{d} - \mathbf{g}(\mathbf{m}) \right] - \lambda \underline{\mathbf{R}} \mathbf{m}_k \right)$$

where at the *k*th iteration, $\underline{\underline{\mathbf{G}}}_{k} = \frac{\partial \underline{\mathbf{g}}(\underline{\mathbf{m}})}{\partial \mathbf{m}}\Big|_{\mathbf{m}=\mathbf{m}_{k}}$, is the appropriate Jacobian matrix, $\underline{\underline{\mathbf{R}}}$ is a regularization matrix discussed in detail below, and $\underline{\underline{\mathbf{W}}}_{k}$ is a (Huber) weighting matrix.

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 $\underline{\mathbf{W}}_{k}$ is are updated at each iteration, and consists of diagonal elements

$${}^{k}w_{i} = \min\left(1, \frac{c\sigma}{{}^{k}d_{i}}\right).$$

 ${}^{k}d_{i}$ is the scalar residual of the *i*th data point using model vector $\mathbf{m}_{k}\Delta F_{i} = \left|\vec{B}_{i,\text{VFM}}\right| - F_{i,\text{ASM}}$ is the scalar residual of the *i*th data point, with $\vec{B}_{i,\text{VFM}}$ calculated using the model parameters from iteration k, and F_{ASM} being the fully calibrated and corrected scalar field measurements from the ASM scalar magnetometer, and

$$\sigma = \sqrt{\frac{\sum_{i} \left(k^{-1} w_i^{-k} d_i\right)^2}{\sum_{i}^{k-1} w_i^2}}$$

being a (robust) estimate of the standard deviation of the residuals at iteration k.

¹⁷⁵ We set c = 2, slightly higher than the value of 1.5 usually chosen, in order to ensure

that the less numerous polar data are not overly downweighted in the determination

- ¹⁷⁷ of the calibration parameters.
- 178

It turns out that the full set of 2,046 parameters is not needed to obtain good re-179 sults and low data misfit, which is confirmed by inspection of the eigenvalues of the 180 matrix $(\underline{\mathbf{G}}_{k}^{T}\underline{\mathbf{W}}_{k}\underline{\mathbf{G}}_{k}+\lambda\underline{\mathbf{R}})$, as presented in Fig. 4 for Swarm Alpha. The magnitudes 181 of the sorted eigenvalues (in order of decreasing magnitude) exhibit a distinct drop 182 around 750-800 degrees of freedom, indicating the smaller eigenvalues contribute 183 little to the solution. The inversion of this matrix was therefore finally performed 184 using a truncated singular value decomposition (TSVD) procedure, retaining only 185 750 degrees of freedom. 186

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A regularization matrix $\underline{\mathbf{R}}$ is also included to help stabilizeto stable the inver-188 siom. This is necessary because the *Swarm* satellites operates in a tightly controlled 189 attitude orientation which leads to a poor excitation of the VFM instrument along 190 the axis perpendicular to the orbit plane (the East-West direction corresponding 191 to the y-axis of the VFM sensor). Consequently, the parameters related to the y-192 axis are poorly determined in a scalar calibration. The regularization matrix $\underline{\mathbf{R}}$ is 193 therefore defined so that it acts on the parameters $s_{2,\text{Tsensor}}$, $s_{2,\beta}$, u_1 , and u_3 to 194 force $s_{2,\text{Tsensor}} \simeq (s_{1,\text{Tsensor}} + s_{3,\text{Tsensor}})/2$ (to reflect the physical properties of the 195 VFM sensor) and also to minimize the norms $s_{2,\beta}^2$ and $u_1^2 + u_3^2$. λ is chosen to be 196 sufficiently large to effectively impose the regularisation on the estimated model. 197 Note that no regularisation is directly imposed on $\delta \vec{B}_{Sun}$ but use of truncated SVD 198 during the inversion automatically acts to suppresses structure in regions that are 199 not well constrained by the input data. 200

201

The starting model for the inversions is "unity", i.e. $\underline{\underline{P}} = \underline{\underline{S}} = \underline{\underline{I}}$, where $\underline{\underline{I}}$ is the identity matrix, and $\vec{u}_n^m = \vec{v}_n^m = \vec{0}$. The inversions typically converge within 204 25 iterations.

²⁰⁵ **3** Results of Model Estimation for Swarm Alpha

The model described above is estimated for Swarm Alpha using data from the begin-206 ning of the mission (22 November 2013) until June 2015. Fig. 1 shows the final scalar 207 residuals, i.e. the residuals after application of the model (after "calibration and 208 correction") of the VFM measurements, (in green) as a function of time together 209 with the residuals of the un-corrected but re-scaled vector field measurements, i.e. 210 $\vec{B}_{\rm VFM} + \delta \vec{B}_{\rm Sun}$, in light blue; these data illustrate what can be achieved with the 211 traditional scalar calibration methods. Note the excellent reduction of the scalar 212 residuals achieved by the model; the Huber weighted rms of the residuals drops 213 from 963 pT to 168 pT. Table 2 provides the corresponding numbers for Bravo and 214 Charlie. 215

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Fig. 5 shows normal distribution plots for the scalar residuals. The top plot shows

the distributions of all data for un-corrected (red) and fully corrected data (green)

and demonstrates a transition from a non-Gaussian to Gaussian residual distributions when applying the model. The bottom plots show the distributions of the
data split into 3-months periods, un-corrected to the left and corrected to the right.
These also demonstrate the elimination of systematic and non-Gaussian effects.

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Table 3 lists the estimated s_{Tsensor} and s_{β} parameters, and the non-orthogonality values for all three *Swarm* satellites together with their estimated pre-flight values for the VFM instrument itself for reference. I.e. the table shows the adjustments applied in order to reduce the scalar residuals to the level indicated above.

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Table 4 shows the increase in the weighted rms of the scalar residuals when omit-229 ting individual parts of the model – a full re-estimation of the remaining model 230 parameters is carried out for each table entry. Particularly the omission of the non-231 orthogonalities drastically increases the misfit – the power (the mean-square) is 232 more than doubled. Due to the stable attitude of the Swarm satellites, the small 233 x-z non-orthogonality angle, u_2 , is equivalent to first order to a small, relative 234 timeshift between the ASM and VFM measurements -1 arc-second corresponds 235 roughly to a 3 ms timeshift, and it has been discussed whether it would be more 236 reasonable to introduce such timeshifts rather than adjusting the pre-flight esti-237 mated non-orthogonalities. However, the variations in the u_2 angles estimated by 238 this model would imply time-shifts varying from -3 ms to +13 ms for the individual239 satellites which, to the authors, seems quite unlikely. 240

241

The temporal evolution of the scaling of the vector field measurements, $s^{B-spline}$, 242 is shown in Fig. 6 for the various test models listed in Table 4. The full model, shown 243 in redgreen, shows a smooth behaviour in time, as expected from an instrument 244 design perspective. The blue curve shows the model without s_{β} ; this exhibits some 245 small oscillations, whereas the light brownred (no s_{Tsensor}) and greenmagneta (no 246 $\delta \vec{B}_{Sun}$) curves show much higher level of oscillations indicating they are inadequate 247 to capture the behaviour of the measurements. The elimination of the oscillations 248 in the full model is a good indicator of the validity of this model. The magentaeyan 249 curve shows the model without non-orthogonalities; this is rather close to the curve 250 of the full model and indicates the decoupling of the non-orthogonalities from any 251 long-term temporal effect of the measurement disturbances and instruments. 252

253

Maps of the three components of the estimated disturbance fields from the full 254 model as function of Sun incidence angles α (abscissa) and β (ordinate) are given 255 in Figs. 7, 8, and 9 for Swarm Alpha, Bravo, and Charlie respectively. During nom-256 inal flight, the Sun incidence angles traverse these plots horizontally from right 257 to left, and move up or down in β as the orbit plane moves through local time. 258 The Sun induced disturbance is observed to have temporal characteristics that are 259 observed in the plots as horizontally stretched features, these are attributed to 260 thermal capacitance: The Sun induced disturbance exhibits characteristic warm-up 26 and cool-down effects, i.e. the disturbance increases when the spacecraft is exposed 262 to the Sun, and decreases when the Sun exposure terminates. The time constants 263 for these effects are up to tens of minutes (corresponding to several tens of degrees 264

²⁶⁵ in the α angle). This effect is captured by the spherical harmonic model expansion ²⁶⁶ of $\delta \vec{B}_{Sun}$ and yields the horizontally stretched features in Figs. 7-9. Note also the ²⁶⁷ regions of nightside data (eclipse), the circled areas to the left of the figures, which ²⁶⁸ generally show less disturbance; this is not imposed by the model or any regulari-²⁶⁹ sation, rather it is simply a result of the data itself, and thus another indicator of ²⁷⁰ the ability of the model to describe the observed disturbances. The plots also show ²⁷¹ both the similarities and the differences in $\delta \vec{B}_{Sun}$ between the three satellites.

273 4 Conclusions

We have established a predominantly empirical model for the calibration and cor-274 rection of the magnetic vector field measurements of the three Swarm spacecraft. 275 The model is based on detailed studies of the observed scalar residuals between the 276 measurements of the absolute scalar magnetometer, ASM, and the modulus of the 277 measurements of the vector field magnetometer, VFM. The model has proven to be 278 quite robust as more data are incorporated into the estimation of the model parame-279 ters, although the ambiguity of determining vector disturbances from a pure scalar 280 calibration affects the estimated correction vectors; these corrections do change 281 slightly (by a few tenths of a nT) as more data are added. 282

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The estimated models reduce the scalar differences between the *Swarm* magnetometers to generally below 0.5 nT with rms values well below 200 pT for all three satellites, and have been in operational use since April 2015 to produce corrected *Swarm* Level 1b magnetic field vector data (as of version 0401).

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Future evolutions of the model presented here are foreseen to include changing the model of the temporal evolution of the VFM sensitivity from B-splines to an exponentially decaying function. Analysis of $\delta \vec{B}_{Sun}$ also indicates that this vector is generally confined to a few, distinct directions which may be incorporated in future models. Finally, it may be possible to model the effect of the thermal capacitance using appropriate temporal filter functions which would lead to a significant reduction of the number of parameters of the model.

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297 Data Availability

The estimated disturbance vectors, $\delta \vec{B}_{Sun}$, are included in the operational Level 1b magnetic *Swarm* data products as dB_Sun.

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³⁰¹ Uncorrected data are available at *ftp://swarm-diss.eo.esa.int/Advanced/* (login ³⁰² required, access can be requested via *https://earth.esa.int/Swarm*).

303

304 Competing interests

306 Author's contributions

The authors declare that they have no competing interests.

³⁰⁷ LTC carried out the in-flight scalar calibration and characterisation, analysed the results, and led the writing of this

manuscript. VL proposed the model for the Sun induced vector disturbance, $\delta B_{
m Sun}$, and made the first estimations

using this model. NiO and CF supported the entire project with many discussions, suggestions, and source code.

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- 317 of the manuscript.
- This paper is the IPGP contribution XXXX. 318

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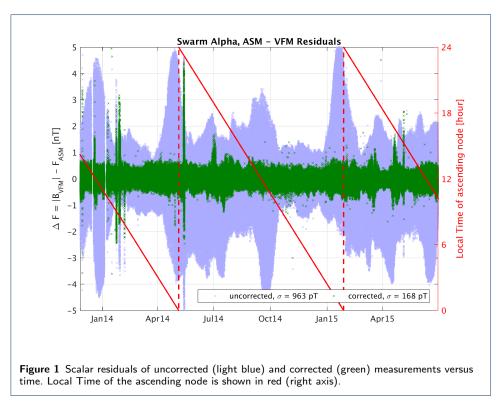
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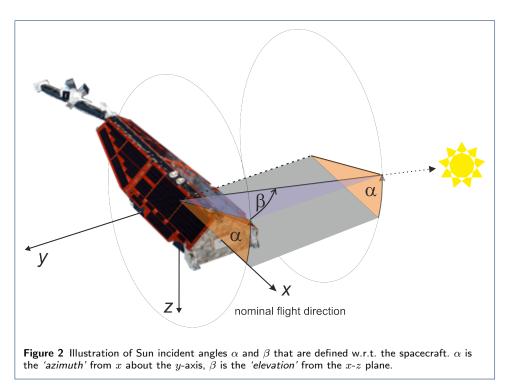
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335 Figures





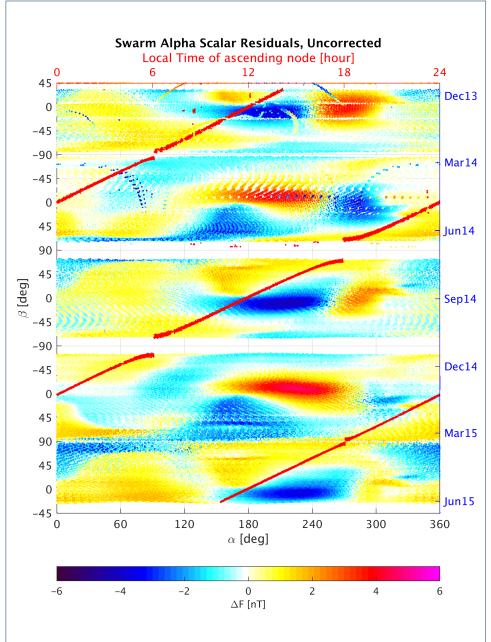
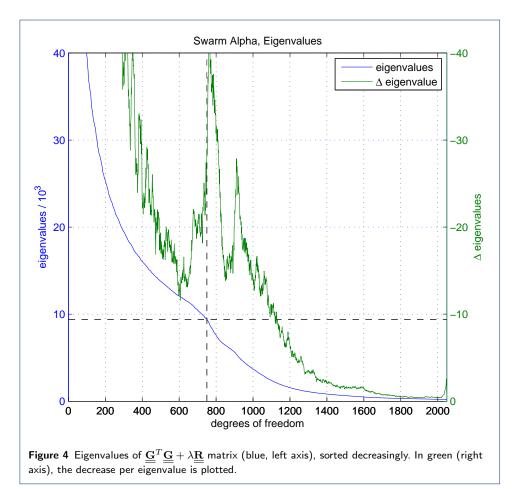
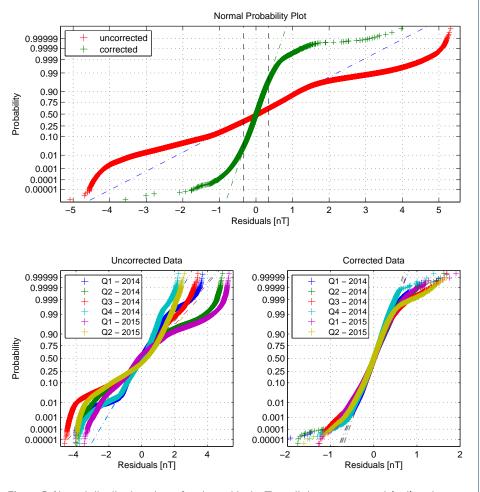
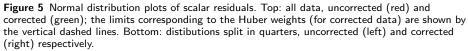


Figure 3 Uncorrected scalar residuals between ASM and VFM magnetometers, ΔF , plotted versus Sun incident angles α and β . The β angle oscillates slowly in time, hence the *un-folded* β angle corresponds to season which is indicated in blue on the right hand side. Local Time of the ascending node is shown in red (axis on top).







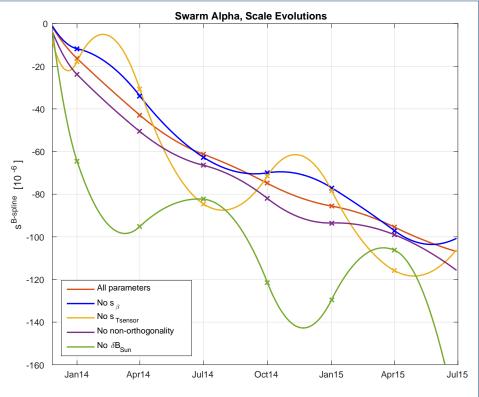
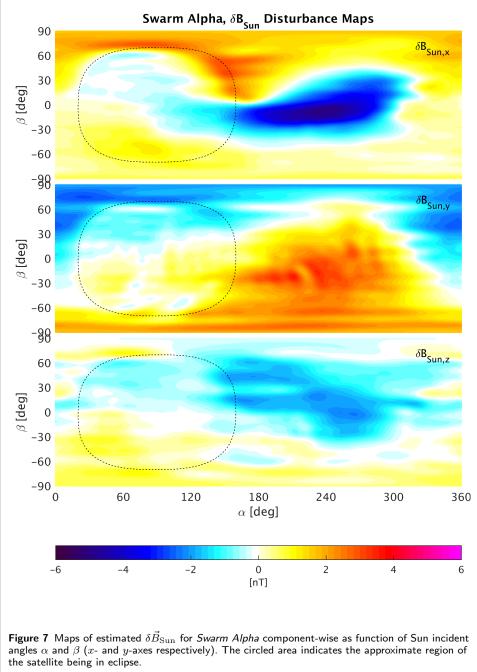
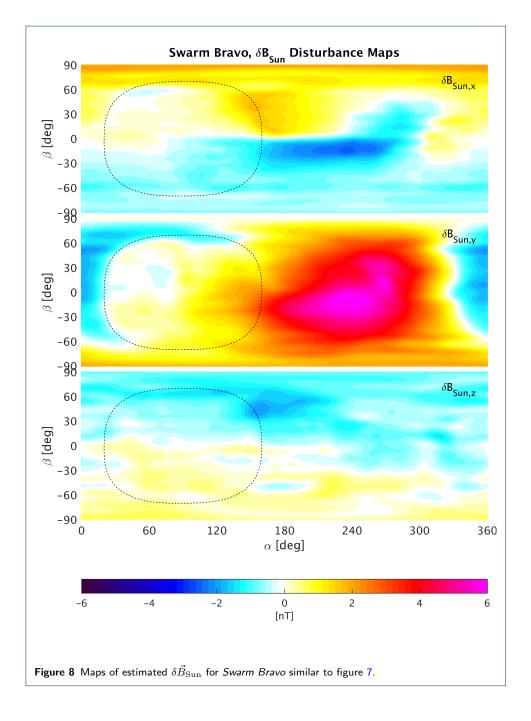
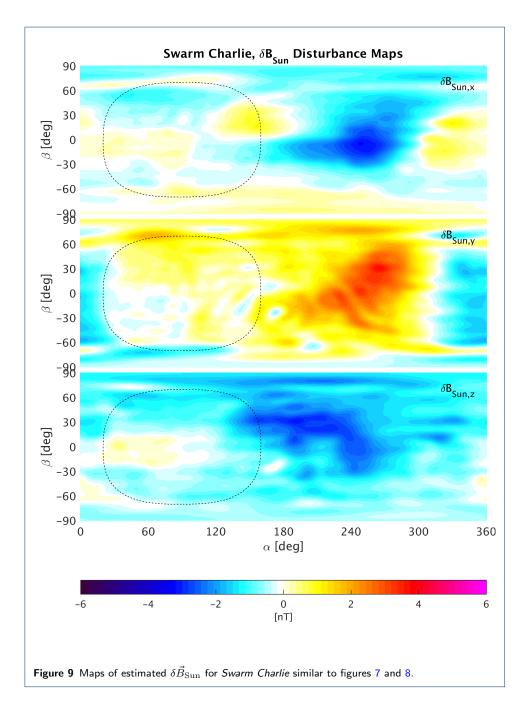


Figure 6 Temporal evolution of estimated VFM scaling for various test models. RedGreen: full model. Blue: full model except s_{β} . Light brownRed: full model, except adjustment of s_{Tsensor} . MagentaCyan: full model, except adjustment of VFM sensor non-orthogonalities. GreenMagenta: full model, except disturbance $(\delta \vec{B}_{\mathrm{Sun}})$. B-spline knots are marked with crosses.







336 Tables

Table 1 Model parameters

Description	Parameters	Dimension
$\delta ec{B}_{ m Sun}$	$ec{u},ec{v}$	2,028
Sensitivity, time dependent	$s^{Bspline}$	9
Sensitivity, eta dependency	$ec{s}_eta$	3
Sensitivity, sensor temperature dependency	$ec{s}_{Tsensor}$	3
Non-orthogonalities	u_1, u_2, u_3	3
Total		2,046
Sensitivity, sensor temperature dependency Non-orthogonalities	$ec{s}_{Tsensor}$	3

 Table 2 Scalar Residual Statistics, Uncorrected and Corrected Data.

For Swarm Charlie two sets of numbers are given; one set for which the ASM was still working ($F_{\rm ASM}$, until 5. November 2014) and one set using the scalar data from Swarm Alpha mapped to the position of Swarm Charlie ($F_{\rm AC,map}$). For data from 1. May 2014 through 5. November 2014 the weighted rms of $F_{\rm ASM} - F_{\rm AC,map}$ is 572.6 pT.

Satellite		Weighted rms [pT]	
		Uncorrected	Corrected
Alpha		962.6	168.3
Bravo		710.3	164.2
Charlie	$F_{\rm ASM}$	632.1	172.3
	$F_{\rm AC,map}$	862.1	527.7

Table 3 Estimated values for selected model parameters for all three *Swarm* satellites. The *nT*-equivalents of the adjustments in a 50,000 *nT* ambient field are: $s_{\rm Tsensor} = 10^{-6}/^{\circ}C \sim 1.25 nT$ (25°*C* temperature swing), $s_{\beta} = 0.1 \times 10^{-6}/{\rm deg} \sim \pm 0.45 nT$ (±90 deg), u = 1 arc-second ~ 0.242 *nT*.

Sat	temperatu	ity/sensor ire, s_{Tsensor} , ${}^{-6}/{}^{\circ}C]$			Non-orthogonalities, $u_{1,2,3}$, [arc-seconds]	
	Pre-flight	Adjustment	Pre-flight	Adjustment	Pre-flight	Adjustment
Alpha	28.5	0.616	_	-0.125	102.386	-0.601
	28.8	0.780		0	217.403	-3.960
	28.3	0.945		0.012	-179.318	0.149
Bravo	28.3	1.168	_	-0.132	350.880	-0.558
	29.0	1.385	_	-0.003	62.432	-2.453
	28.8	1.602		-0.198	-147.060	1.608
Charlie	27.7	1.521	-	-0.090	139.140	0.094
	29.1	1.300	-	-0.038	-248.890	1.042
	28.4	1.076		-0.167	-109.960	0.805

Table 4 Weighted rms values for various models, Swarm Alpha

Model	weighted rms [pT]	Residual power (normalized)
Full model	168.3	100%
No s_{β}	176.1	107%
No $s_{\rm Tsensor}$	181.7	116%
No non-orthogonalities	250.2	221%
No $\delta ec{B}_{ m Sun}$	962.6	3,269%